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A General and Comprehensive Subclass of Univalent Functions Associated with Certain Geometric Functions

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Abstract: In this paper, taking into account the intriguing recent results of Rabotnov functions, Poisson functions, Bessel functions and Wright functions, we consider a new comprehensive subclass $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$ of univalent functions defined in the unit disk $\Lambda = \{\tau \in \mathbb{C} : |\tau| < 1\}$. More specifically, we investigate some sufficient conditions for Rabotnov functions, Poisson functions, Bessel functions and Wright functions to be in this subclass. Some corollaries of our main results are given. The novelty and the advantage of this research could inspire researchers of further studies to find new sufficient conditions to be in the subclass $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$ not only for the aforementioned special functions but for different types of special functions, especially for hypergeometric functions, Dini functions, Sturve functions and others.

Keywords: analytic functions; Rabotnov function; Poisson distribution; Bessel function; Wright function

MSC: 30C45



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1. Introduction and Preliminaries

In mathematics, hypergeometric functions are a basic family of special functions with wide-ranging applications in number theory, probability theory, mathematical physics, combinatorics and differential equations. In addition to serving as building blocks for creating mathematical models and solving issues in a variety of scientific and engineering disciplines, they offer solutions to differential equations that describe physical phenomena. Their applicability makes them essential tools in mathematical analysis and scientific research. These functions show many mathematical features and linkages as recurrence relations and solutions to differential equations.

Specific instances of hypergeometric functions can be used to represent a wide range of classical special functions, including the Gaussian hypergeometric function, confluent hypergeometric functions and Legendre functions. Moreover, many correlations exist between hypergeometric functions and other mathematical functions, including gamma and Bessel functions.

In particular, in this paper we will consider Rabotnov functions, Poisson functions, Bessel functions and Wright functions, respectively.

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At first, let \digamma the class of analytic and univalent functions $\mathcal L$ defined in the unit disk $\Lambda = \{\tau \in \mathbb C : |\tau| < 1\}$ and meet $\mathcal L'(0) - 1 = \mathcal L(0) = 0$. Thus, each $\mathcal L \in \digamma$ has the following series:

$$\mathcal{L}(\tau) = \tau + a_2 \tau^2 + a_3 \tau^3 + \dots = \tau + \sum_{u=2}^{\infty} a_u \tau^u, \ (\tau \in \Lambda).$$
 (1)

The Rabotnov function as a special function was introduced by Rabotnov [1] in 1948, as follows:

$$\Phi_{\ell,\beta}(\tau) = \tau^{\ell} \sum_{u=0}^{\infty} \frac{(\beta)^{u} \tau^{u(\ell+1)}}{\Gamma((u+1)(\ell+1))}, \quad (\ell,\beta,\tau \in \mathbb{C}).$$

The well-known Mittag-Leffler function, which is frequently used to solve integral equations, is represented by the Rabotnov function in particular cases.

The connection between the Mittag–Leffler function M and the Rabotnov function $\Phi_{\ell,\beta}(\tau)$ is shown in the recurrence relation

$$\Phi_{\ell,\beta}(\tau) = \tau^{\ell} M_{\ell+1,\ell+1}(\beta \tau^{\ell+1}), \ \ell,\beta,\tau \in \mathbb{C}.$$

For additional details regarding Mittag-Leffler, see [2-5].

The Rabotnov function $\Phi_{\ell,\beta}(\tau)$ is evidently not in class F. Consequently, it makes sense to take into account the normalizing of the Rabotnov function that follows:

$$\mathcal{R}_{\ell,\beta}(\tau) = \tau^{\frac{1}{\ell+1}} \Gamma(\ell+1) \Phi_{\ell,\beta}(\tau^{\frac{1}{\ell+1}}) = \tau + \sum_{u=2}^{\infty} \frac{\beta^{u-1} \Gamma(\ell+1)}{\Gamma((\ell+1)u)} \tau^{u}, \ \tau \in \Lambda$$
 (2)

where $\ell > -1$ and $\beta \in \mathbb{C}$.

A variable Y is considered Poisson-distributed if the values are taken to be $0,1,2,3,\cdots$ for the parameter ν with probabilities $e^{-\nu}$, $\nu^2 \frac{e^{-\nu}}{1!}$, $\nu^2 \frac{e^{-\nu}}{2!}$, $\nu^3 \frac{e^{-\nu}}{3!}$, \cdots , respectively. Thus,

$$P(Y = s) = \frac{v^s e^{-v}}{s!}, s = 0, 1, 2, 3, \cdots$$

In 2014, the Poisson function was introduced by Porwal [6] (see also [7,8]), as follows:

$$\mathcal{K}(\nu,\tau) = \tau + \sum_{u=2}^{\infty} \frac{\nu^{u-1}}{(u-1)!} e^{-\nu} \tau^{u}, \qquad (\nu > 0, \tau \in \Lambda).$$
 (3)

The radius of convergence for $\mathcal{K}(\nu, \tau)$ is infinite.

The generalized Bessel function of the first kind of $\chi(\tau)$, mentioned as:

$$\chi(\tau) = \chi_{p,b,c}(\tau) = \sum_{u=0}^{\infty} \frac{(-1)^{u} c^{u}}{u! \Gamma(p + \frac{b+1}{2} + u)} \left(\frac{\tau}{2}\right)^{2u+p}, \ \tau \in \mathbb{C}$$
 (4)

is one specific solution to the linear differential equation of second order

$$\tau^{2}\chi''(\tau) + b\tau\chi'(\tau) + [c\tau^{2} - p^{2} + (1-b)p]\chi(\tau) = 0, \quad p, b, c \in \mathbb{C}.$$
 (5)

The function $\chi(\tau)$ is not univalent in symmetric domain Λ . Now, we examine the function $B_{v,b,c}$ as follows:

$$B_{p,b,c}(\tau) = \epsilon_p = 2^p \Gamma(p + \frac{b+1}{2}) \tau_{p,b,c}^{\frac{-p}{2}} \chi(\tau^{\frac{1}{2}}).$$

Using the Pochhammer symbol, defined for $a \neq 0, -1, -2, \cdots$ by

$$(\zeta)_u = \frac{\Gamma(\zeta+u)}{\Gamma(\zeta)} = \left\{ \begin{array}{ll} 1, & \text{if } u=0 \text{ and } \zeta \in \mathbb{C} - \{0\}; \\ \zeta(\zeta+1)...(\zeta+u-1), & \text{if } u \in \mathbb{N} \text{ and } \zeta \in \mathbb{C} \end{array} \right.$$

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we can rewrite the function $B_{v,b,c}$ as follows:

$$B_{p,b,c}(\tau) = \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^u}{\left(p + \frac{b+1}{2}\right)_u} \frac{\tau^u}{u!}$$

and define the function $\mathcal{L}_{p,b}(u,c;\tau)$ as

$$\mathcal{L}_{p,b}(u,c;\tau) = \tau B_{p,b,c}(\tau) = \tau + \sum_{u=2}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-1)!} \tau^u \equiv \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^u}{(\delta)_u u!} \tau^u \tag{6}$$

where $c < 0, -1, -2, \cdots$ and $\delta = p + \frac{b+1}{2} \neq 0$. The function $\mathcal{L}_{p,b}(u, c; \tau) = \mathcal{L}_p(\tau)$ is entire and satisfies the equation

$$4\tau^{2}\mathcal{L}_{p}''(\tau) + 2(2p + b + 1)\tau\mathcal{L}_{p}'(\tau) + c\tau\mathcal{L}_{p}(\tau) = 0.$$
 (7)

The Bessel function is essential in many areas of mathematical physics and applied mathematics, including as signal processing, hydrodynamics, radio physics and acoustics. Thus, a great deal of research has been conducted on Bessel functions. For example, Baricz et al. [9] gave sufficient conditions for Bessel functions. Frasin et al. in [10] determined some conditions for the function $z(2-u_p(z))$, where u_p denotes the Bessel function of order p to be in various subclasses of analytic functions. Saiful et al. in [11] determined various conditions in which Bessel functions have certain geometric properties in the unit disc.

Finally, we will examine the Wright special function, which is defined as

$$\sigma(\varphi, \varepsilon; \tau) = \sum_{u=0}^{\infty} \frac{1}{\Gamma(\varphi u + \varepsilon)} \frac{\tau^{u}}{u!}, \varphi > -1, \varepsilon, \tau \in \mathbb{C}.$$
 (8)

Wright functions have been mentioned in papers pertaining to partial differential equations and in other applications; see, for example [12–15].

Remark 1. For $\varphi = 1$, $\varepsilon = p + 1$ and the functions $\sigma(1, p + 1; -\tau^2/4)$, it is possible to write the following using the Bessel functions $J_p(\tau)$:

$$J_p(\tau) = \left(\frac{\tau}{2}\right)^p \sigma(1, p+1; \frac{-\tau^2}{4}) = \sum_{u=0}^{\infty} \frac{1}{\Gamma(u+p+1)} \frac{(\tau/2)^{2u+p}}{u!}.$$

Furthermore, for $\varphi > 0$ and p > -1, the function $\sigma(\varphi, p + 1; -\tau) \equiv J_p^{\varphi}(\tau)$ is generalized Bessel function.

Observe that the Wright function $\sigma(\varphi, \varepsilon, \tau)$ is not in class F. Thus, we define the next two Wright functions:

$$\mathcal{W}^{(1)}(\varphi, \varepsilon; \tau) := \Gamma(\varepsilon)\tau\sigma(\varphi, \varepsilon; \tau) = \sum_{u=0}^{\infty} \frac{\Gamma(\varepsilon)}{\Gamma(\varphi u + \varepsilon)} \frac{\tau^{u+1}}{u!}$$
$$= \tau + \sum_{u=2}^{\infty} \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1) + \varepsilon)} \frac{\tau^{u}}{(u-1)!}, \quad \varphi > -1, \, \varepsilon > 0, \, \tau \in \Lambda$$
(9)

and

$$\mathcal{W}^{(2)}(\varphi, \varepsilon; \tau) := \Gamma(\varphi + \varepsilon) \left(\sigma(\varphi, \varepsilon; \tau) - \frac{1}{\Gamma(\varepsilon)} \right) = \sum_{u=0}^{\infty} \frac{\Gamma(\varphi + \varepsilon)}{\Gamma(\varphi u + \varphi + \varepsilon)} \frac{\tau^{u+1}}{(u+1)!}$$

$$= \tau + \sum_{u=2}^{\infty} \frac{\Gamma(\varphi + \varepsilon)}{\Gamma(\varphi(u-1) + \varphi + \varepsilon)} \frac{\tau^{u}}{u!}, \quad \varphi > -1, \ \varphi + \varepsilon > 0, \ \tau \in \Lambda.$$
 (10)

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Furthermore, observe that $W^{(1)}(\varphi, \varepsilon; \tau)$, $W^{(2)}(\varphi, \varepsilon; \tau)$ and $H_{\varphi, \varepsilon}(\tau) = \frac{W^{(1)}(\varphi, \varepsilon; \tau)}{\tau}$ are satisfied by the following relations, respectively

$$\begin{split} \varphi\tau(\mathcal{W}^{(1)}(\varphi,\varepsilon;\tau))' &= (\varepsilon-1)\mathcal{W}^{(1)}(\varphi,\varepsilon-1;\tau) + (\varphi+\varepsilon+1)\mathcal{W}^{(1)}(\varphi,\varepsilon;\tau), \\ \tau(\mathcal{W}^{(2)}(\varphi,\varepsilon;\tau))' &= \mathcal{W}^{(1)}(\varphi,\varphi+\varepsilon;\tau) \end{split}$$

and

$$H'_{\varphi,\varepsilon}(\tau) = \frac{\Gamma(\varepsilon)}{\Gamma(\varphi+\varepsilon)} H_{\varphi,\varphi+\varepsilon}(\tau).$$

Recently, numerous scholars have examined classes of analytic functions that involve special functions in order to determine certain conditions in Λ . Many features, generalizations and applications of many kinds of geometric functions have been covered widely in the literature. For example, Miller et al. in [16] determined conditions for the univalence of Gaussian. Ponnusamy et al. in [17] found conditions for function f to be starlike or convex. Ponnusamy et al. in [18] studed sufficient and necessary conditions, in terms of the coefficient A_n , for a function $f \in \mathcal{A}$ to be in subclasses of univalent functions. Ponnusamy et al. in [19] determined conditions for convexity and starlikeness. Yagmur et al. in [20] presented some applications of convexity involving Struve functions. In [21], Frasin et al. provided conditions for the Struve functions to be in two classes of analytic functions.

In this paper, we introduce a new subclass of analytic functions. $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$ of the class \digamma , which generalizes many of the previous classes of analytic functions defined in Λ .

Definition 1. A function $\mathcal{L} \in \mathcal{F}$ is said to be in class $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$ if

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) |a_u| \le \mu,$$

where Δ_1 , Δ_2 , Δ_3 and Δ_4 are real numbers and $\mu > 0$.

By precisely specializing for the coefficients Δ_1 , Δ_2 , Δ_3 , Δ_4 and μ in Definition 1, we note the following:

(1) A function $\mathcal{L} \in \wp(\lambda, \gamma) \equiv \mathcal{O}_{1-\gamma}(0, \lambda, 1-\lambda, -\gamma)$ (which is due to [22]) if

$$\sum_{u=2}^{\infty} (u + \lambda u(u-1) - \gamma)|a_u| \le 1 - \gamma.$$

(2) A function $\mathcal{L} \in \mathcal{K}(\lambda, \gamma) \equiv \mathcal{O}_{1-\gamma}(\lambda, 1-\lambda, -\gamma, 0)$ (which is due to [23]) if

$$\sum_{u=2}^{\infty} u(u + \lambda u(u-1) - \gamma)|a_u| \le 1 - \gamma.$$

(3) A function $\mathcal{L} \in \wp_{\lambda}^*(\gamma, \beta) \equiv \mathcal{O}_{1-\gamma}(0, 0, \beta+1, -\lambda(\gamma+\beta))$ (which is due to [24]) if

$$\sum_{u=2}^{\infty} (u(\beta+1) - \lambda(\gamma+\beta))|a_u| \le 1 - \gamma.$$

(4) A function $\mathcal{L} \in \bigoplus_{\lambda}^* (\gamma, \beta) \equiv \mathcal{O}_{1-\gamma}(0, \beta+1, -\lambda(\gamma+\beta), 0)$ (which is due to [24]) if

$$\sum_{u=2}^{\infty} u(u(\beta+1) - \lambda(\gamma+\beta))|a_u| \le 1 - \gamma.$$

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(5) A function $\mathcal{L} \in \mathcal{T}(\lambda, \alpha) \equiv \mathcal{O}_{1-\alpha}(0, 0, 1-\alpha\lambda, -\alpha(1-\lambda))$ (which is due to [25]) if

$$\sum_{u=2}^{\infty} (u(1-\alpha\lambda) - \alpha(1-\lambda))|a_u| \le 1-\alpha.$$

(6) A function $\mathcal{L} \in \mathcal{C}(\lambda, \alpha) \equiv \mathcal{O}_{1-\alpha}(0, 1-\alpha\lambda, -\alpha(1-\lambda), 0)$ (which is due to [25]) if

$$\sum_{u=2}^{\infty} u(u(1-\alpha\lambda)-\alpha(1-\lambda))|a_u| \leq 1-\alpha.$$

(7) A function $\mathcal{L} \in \mathcal{SP}_v(\sigma, v) \equiv \mathcal{O}_{\cos \sigma - v}(0, 0, 2, -\cos \sigma - v)$ (which is due to [26]) if

$$\sum_{u=2}^{\infty} (2u - \cos \sigma - v)|a_u| \le \cos \sigma - v.$$

(8) A function $\mathcal{L} \in \mathcal{UCSP}_p(\sigma, v) \equiv \mathcal{O}_{\cos \sigma - v}(0, 2, -\cos \sigma - v, 0)$ (which is due to [26]) if

$$\sum_{u=2}^{\infty} u(2u - \cos \sigma - v)|a_u| \le \cos \sigma - v.$$

To give sufficient conditions for the Rabotnov function, Poisson distribution function, Bessel function and Wright functions to be in the comprehensive subclass $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$, we need the next Lemma given by Sümer Eker [27]:

Lemma 1. *For* $u \in \mathbb{N}$ *and* $\hbar \geq 0$ *, then*

$$(\hbar + 1)^{u-1}(u-1)!\Gamma(\hbar + 1) \le \Gamma(u(\hbar + 1)).$$

Furthermore, from this Lemma, we can write

$$\frac{\Gamma(\hbar+1)}{\Gamma(u(\hbar+1))} \le \frac{1}{(\hbar+1)^{u-1}(u-1)!}.$$
(11)

Furthermore, since $\Gamma(u-1+\beth) \leq \Gamma(\varphi(u-1)+\beth)$ for $\beth>0$ and $u\in\mathbb{C}$, the Lemma holds. We can write

$$\frac{\Gamma(\beth)}{\Gamma(\varphi(u-1)+\beth)} \le \frac{1}{(\beth)_{u-1}}, \quad u \in \mathbb{C}$$
 (12)

Further, since

$$(\beth)_{u-1} = \beth(\beth+1)(\beth+2)\cdots(\beth+u-2) \ge \beth(\beth+1)^{u-2}, \quad u \in \mathbb{C}$$

and using (12), then

$$\frac{\Gamma(\beth)}{\Gamma(\varphi(u-1)+\beth)} \le \frac{1}{\beth(\beth+1)^{u-2}}.$$
(13)

2. Main Results

In this part, we examine a few prerequisites that must be met for Rabotnov functions, Poisson functions, Bessel functions and Wright functions to be in the comprehensive subclass $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$.

Theorem 1. *If the following inequality is valid,*

$$\left(\frac{\Delta_{1}\beta^{3}}{(\ell+1)^{3}} + \frac{(6\Delta_{1} + \Delta_{2})\beta^{2}}{(\ell+1)^{2}} + \frac{(7\Delta_{1} + 3\Delta_{2} + \Delta_{3})\beta}{\ell+1} + \Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}\right)e^{\frac{\beta}{\ell+1}} \\
< \mu + \Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}.$$

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for $\mu > 0$, then the Rabotnov function $\mathcal{R}_{\ell,\beta}(\tau)$ is in the class $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$.

Proof. By Definition 1 and the Rabotnov function given by (2), it suffices to show that

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\beta^{u-1} \Gamma(\ell+1)}{\Gamma((\ell+1)u)} \leq \mu.$$

Let

$$\hbar_1(\lambda,\beta;\mu) = \sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4\right) \frac{\beta^{u-1} \Gamma(\ell+1)}{\Gamma((\ell+1)u)}.$$

Setting

$$\begin{cases}
 u = (u-1) + 1; \\
 u^2 = (u-1)(u-2) + 3(u-1) + 1; \\
 u^3 = (u-1)(u-2)(u-3) + 6(u-1)(u-2) + 7(u-1) + 1,
\end{cases} (14)$$

we obtain

$$\begin{split} \hbar_1(\lambda,\beta;\mu) &= \sum_{u=2}^\infty \left\{ \Delta_1(u-1)(u-2)(u-3) + (6\Delta_1 + \Delta_2)(u-1)(u-2) \right. \\ &+ (7\Delta_1 + 3\Delta_2 + \Delta_3)(u-1) + \Delta_1 + \Delta_2 + \Delta_3 + \Delta_4 \right\} \frac{\beta^{u-1}\Gamma(\ell+1)}{\Gamma((\ell+1)u)}. \\ &= \sum_{u=2}^\infty \frac{\Delta_1(u-1)(u-2)(u-3)\beta^{u-1}\Gamma(\ell+1)}{\Gamma((\ell+1)u)} \\ &+ \sum_{u=2}^\infty \frac{(6\Delta_1 + \Delta_2)(u-1)(u-2)\beta^{u-1}\Gamma(\ell+1)}{\Gamma((\ell+1)u)} \\ &+ \sum_{u=2}^\infty \frac{(7\Delta_1 + 3\Delta_2 + \Delta_3)(u-1)\beta^{u-1}\Gamma(\ell+1)}{\Gamma((\ell+1)u)} \\ &+ \sum_{u=2}^\infty \frac{(\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4)\beta^{u-1}\Gamma(\ell+1)}{\Gamma((\ell+1)u)}. \end{split}$$

Under hypothesis (11), we have

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$$\begin{split} \hbar_{1}(\lambda,\beta;\mu) &\leq \Delta_{1} \sum_{u=2}^{\infty} \frac{(u-1)(u-2)(u-3)\beta^{u-1}}{(\ell+1)^{u-1}(u-1)!} + (6\Delta_{1} + \Delta_{2}) \sum_{u=2}^{\infty} \frac{(u-1)(u-2)\beta^{u-1}}{(\ell+1)^{u-1}(u-1)!} \\ &+ (7\Delta_{3} + 3\Delta_{2} + \Delta_{3}) \sum_{u=2}^{\infty} \frac{(u-1)\beta^{u-1}}{(\ell+1)^{u-1}(u-1)!} \\ &+ (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \sum_{u=2}^{\infty} \frac{\beta^{u-1}}{(\ell+1)^{u-1}(u-1)!} \\ &= \Delta_{1} \sum_{u=4}^{\infty} \frac{\beta^{u-1}}{(\ell+1)^{u-1}(u-4)!} + (6\Delta_{1} + \Delta_{2}) \sum_{u=3}^{\infty} \frac{\beta^{u-1}}{(\ell+1)^{u-1}(u-3)!} \\ &+ (7\Delta_{1} + 3\Delta_{2} + \Delta_{3}) \sum_{u=2}^{\infty} \frac{\beta^{u-1}}{(\ell+1)^{u-1}(u-2)!} \\ &+ (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \sum_{u=2}^{\infty} \frac{\beta^{u-1}}{(\ell+1)^{u-1}(u-1)!} \\ &= \frac{\Delta_{1}\beta^{3}}{(\ell+1)^{3}} e^{\frac{\beta}{\ell+1}} + \frac{(6\Delta_{1} + \Delta_{2})\beta^{2}}{(\ell+1)^{2}} e^{\frac{\beta}{\ell+1}} + \frac{(7\Delta_{1} + 3\Delta_{2} + \Delta_{3})\beta}{\ell+1} e^{\frac{\beta}{\ell+1}} \\ &+ (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \left(e^{\frac{\beta}{\ell+1}} - 1\right). \end{split}$$

The right-hand term in (15) is bounded above by μ ; thus,

$$\left(\frac{\Delta_{1}\beta^{3}}{(\ell+1)^{3}} + \frac{(6\Delta_{1} + \Delta_{2})\beta^{2}}{(\ell+1)^{2}} + \frac{(7\Delta_{1} + 3\Delta_{2} + \Delta_{3})\beta}{\ell+1} + \Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}\right)e^{\frac{\beta}{\ell+1}}$$

$$\leq \mu + \Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}.$$

This concludes the proof of Theorem 1. \Box

Theorem 2. If the following inequality is valid,

$$\Delta_1 \nu^3 + (6\Delta_1 + \Delta_2)\nu^2 + (7\Delta_1 + 3\Delta_2 + \Delta_3)\nu + (\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4)(1 - e^{-\nu}) \le \mu.$$

for $\mu > 0$, then the Poisson distribution function $\mathcal{K}(\nu, \tau)$ belongs to $\mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$.

Proof. In view of Definition 1 and and Equality (3) for the Poisson function, it suffices to show that

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{v^{u-1}}{(u-1)!} e^{-v} \le \mu.$$

Let

$$h_2(\nu,\mu) = \sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\nu^{u-1}}{(u-1)!} e^{-\nu}.$$

Similar to Theorem 1, we obtain

$$\begin{split} \hbar_2(\nu,\mu) &= \Delta_1 \sum_{u=4}^{\infty} \frac{\nu^{u-1}}{(u-4)!} e^{-\nu} + (6\Delta_1 + \Delta_2) \sum_{u=3}^{\infty} \frac{\nu^{u-1}}{(u-3)!} e^{-\nu} \\ &+ (7\Delta_1 + 3\Delta_2 + \Delta_3) \sum_{u=2}^{\infty} \frac{\nu^{u-1}}{(u-2)!} e^{-\nu} + (\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4) \sum_{u=2}^{\infty} \frac{\nu^{u-1}}{(u-1)!} e^{-\nu}. \\ &= \Delta_1 \nu^3 + (6\Delta_1 + \Delta_2) \nu^2 + (7\Delta_1 + 3\Delta_2 + \Delta_3) \nu + (\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4) (1 - e^{-\nu}) \le \mu. \end{split}$$

This concludes the proof of Theorem 2. \Box

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Theorem 3. *If the following inequality is valid,*

$$\Delta_{1} \mathcal{L}_{p}^{""}(1) + (6\Delta_{1} + \Delta_{2}) \mathcal{L}_{p}^{"}(1) + (7\Delta_{1} + 3\Delta_{2} + \Delta_{3}) \mathcal{L}_{p}^{\prime}(1) + (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \mathcal{L}_{p}(1)
\leq \mu + \Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}.$$

for $\mu > 0$, then the Bessel function $\mathcal{L}_{p,b}(u,c;\tau)$ belongs to $\mathcal{O}_{\mu}(\Delta_1,\Delta_2,\Delta_3,\Delta_4)$.

Proof. By Definition 1 and the Bessel function given by (6), it suffices to show that

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-1)!} \le \mu.$$

Let

$$\hbar_3(c,\delta;\mu) = \sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-1)!}.$$

Similar to Theorem 1, we obtain

$$\begin{split} &\hbar_{3}(c,\delta;\mu) = \Delta_{1} \sum_{u=4}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-4)!} + (6\Delta_{1} + \Delta_{2}) \sum_{u=3}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-3)!} \\ &+ (7\Delta_{1} + 3\Delta_{2} + \Delta_{3}) \sum_{u=2}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-2)!} + (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \sum_{u=2}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u-1}}{(\delta)_{u-1}(u-1)!} \\ &= \frac{\Delta_{1}\left(\frac{-c}{4}\right)^{3}}{\delta(\delta+1)(\delta+2)} \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u}}{(\delta+3)_{u-2}u!} + \frac{\left(6\Delta_{1} + \Delta_{2}\right)\left(\frac{-c}{4}\right)^{2}}{\delta(\delta+1)} \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u}}{(\delta+2)_{u-1}u!} \\ &+ \frac{\left(7\Delta_{1} + 3\Delta_{2} + \Delta_{3}\right)\left(\frac{-c}{4}\right)}{\delta} \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u}}{(\delta+1)_{u}u!} + (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \sum_{u=0}^{\infty} \frac{\left(\frac{-c}{4}\right)^{u+1}}{(\delta)_{u+1}(u+1)!} \\ &= \frac{\Delta_{1}\left(\frac{-c}{4}\right)^{3}}{\delta(\delta+1)(\delta+2)} \mathcal{L}_{p+3}(1) + \frac{\left(6\Delta_{1} + \Delta_{2}\right)\left(\frac{-c}{4}\right)^{2}}{\delta(\delta+1)} \mathcal{L}_{p+2}(1) \\ &+ \frac{\left(7\Delta_{1} + 3\Delta_{2} + \Delta_{3}\right)\left(\frac{-c}{4}\right)}{\delta} \mathcal{L}_{p+1}(1) + (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \left(\mathcal{L}_{p}(1) - 1\right) \\ &= \Delta_{1}\mathcal{L}_{p}^{"'}(1) + \left(6\Delta_{1} + \Delta_{2}\right)\mathcal{L}_{p}^{"}(1) + \left(7\Delta_{1} + 3\Delta_{2} + \Delta_{3}\right)\mathcal{L}_{p}^{'}(1) \\ &+ \left(\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}\right) \left(\mathcal{L}_{p}(1) - 1\right) \\ &\leq \mu. \end{split}$$

Theorem 4. If the following inequality is valid,

$$\begin{split} &e^{\frac{1}{\varepsilon+1}} \left(\frac{\Delta_1}{\varepsilon(\varepsilon+1)^2} + \frac{6\Delta_1 + \Delta_2}{\varepsilon(\varepsilon+1)} + \frac{7\Delta_1 + 3\Delta_2 + \Delta_3}{\varepsilon} + \frac{(\varepsilon+1)(\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4)}{\varepsilon} \right) \\ &\leq \mu + \frac{(\varepsilon+1)(\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4)}{\varepsilon} \end{split}$$

for $\mu > 0$, then $\mathcal{W}^{(1)}(\varphi, \varepsilon; \tau) \in \mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$.

Proof. By Definition 1 and the Wright function given by (9), it suffices to show that

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1) + \varepsilon)(u-1)!} \le \mu.$$

Let

$$\hbar_4(\varepsilon,\varphi;\mu) = \sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1) + \varepsilon)(u-1)!}.$$

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By (14), we obtain

$$\begin{split} \hbar_4(\varepsilon,\varphi;\mu) &= \Delta_1 \sum_{u=4}^\infty \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1)+\varepsilon)(u-4)!} + (6\Delta_1 + \Delta_2) \sum_{u=3}^\infty \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1)+\varepsilon)(u-3)!} \\ &+ (7\Delta_1 + 3\Delta_2 + \Delta_3) \sum_{u=2}^\infty \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1)+\varepsilon)(u-2)!} \\ &+ (\Delta_1 + \Delta_2 + \Delta_3 + \Delta_4) \sum_{u=2}^\infty \frac{\Gamma(\varepsilon)}{\Gamma(\varphi(u-1)+\varepsilon)(u-1)!}. \end{split}$$

Under hypothesis (13), we obtain

$$\begin{split} \hbar_{4}(\varepsilon,\varphi;\mu) & \leq \Delta_{1} \sum_{u=4}^{\infty} \frac{1}{\varepsilon(\varepsilon+1)^{u-2}(u-4)!} + (6\Delta_{1} + \Delta_{2}) \sum_{u=3}^{\infty} \frac{1}{\varepsilon(\varepsilon+1)^{u-2}(u-3)!} \\ & + (7\Delta_{1} + 3\Delta_{2} + \Delta_{3}) \sum_{u=2}^{\infty} \frac{1}{\varepsilon(\varepsilon+1)^{u-2}(u-2)!} \\ & + (\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4}) \sum_{u=2}^{\infty} \frac{1}{\varepsilon(\varepsilon+1)^{u-2}(u-1)!} \\ & = \frac{\Delta_{1}}{\varepsilon(\varepsilon+1)^{2}} e^{\frac{1}{\varepsilon+1}} + \frac{6\Delta_{1} + \Delta_{2}}{\varepsilon(\varepsilon+1)} e^{\frac{1}{\varepsilon+1}} + \frac{7\Delta_{1} + 3\Delta_{2} + \Delta_{3}}{\varepsilon} e^{\frac{1}{\varepsilon+1}} \\ & + \frac{(\varepsilon+1)(\Delta_{1} + \Delta_{2} + \Delta_{3} + \Delta_{4})}{\varepsilon} \left(e^{\frac{1}{\varepsilon+1}} - 1\right) \leq \mu. \end{split}$$

Theorem 5. If the following inequality is valid,

$$\begin{split} & \Delta_1(\varphi+\varepsilon+1)e^{\frac{1}{\varphi+\varepsilon+1}} + (3\Delta_1+\Delta_2)e^{\frac{1}{\varphi+\varepsilon+1}} + (\Delta_1+\Delta_2+\Delta_3)(\varphi+\varepsilon+1) \left(e^{\frac{1}{\varphi+\varepsilon+1}}-1\right) \\ & + \Delta_4(\varphi+\varepsilon+1)^2 \left(e^{\frac{1}{\varphi+\varepsilon+1}} - \frac{1}{\varphi+\varepsilon+1}-1\right) \leq (\varphi+\varepsilon)\mu \end{split}$$

for $\mu > 0$, then $\mathcal{W}^{(2)}(\varphi, \varepsilon; \tau) \in \mathcal{O}_{\mu}(\Delta_1, \Delta_2, \Delta_3, \Delta_4)$.

Proof. By Definition 1 and the Wright function given by (10), it suffices to show that

$$\sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\Gamma(\varphi + \varepsilon)}{\Gamma(\varphi(u - 1) + \varphi + \varepsilon) u!} \le \mu.$$

Let

$$\hbar_5(\varepsilon,\varphi;\mu) = \sum_{u=2}^{\infty} \left(\Delta_1 u^3 + \Delta_2 u^2 + \Delta_3 u + \Delta_4 \right) \frac{\Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1) + \varphi + \varepsilon)u!}$$

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Similar to Theorem 1, we obtain:

$$\begin{split} \hbar_5(\varepsilon,\varphi;\mu) &= \sum_{u=2}^\infty \frac{\Delta_1 u^2 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} + \sum_{u=2}^\infty \frac{\Delta_2 u \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} \\ &+ \sum_{u=2}^\infty \frac{\Delta_3 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} + \sum_{u=2}^\infty \frac{\Delta_4 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)u!} \\ &= \sum_{u=2}^\infty \frac{\Delta_1 (u-1)(u-2) \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} + \sum_{u=2}^\infty \frac{(3\Delta_1+\Delta_2)(u-1) \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} \\ &+ \sum_{u=2}^\infty \frac{(\Delta_1+\Delta_2+\Delta_3) \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} + \sum_{u=2}^\infty \frac{\Delta_4 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)u!} \\ &= \sum_{u=3}^\infty \frac{\Delta_1 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-3)!} + \sum_{u=2}^\infty \frac{(3\Delta_1+\Delta_2) \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-2)!} \\ &+ \sum_{u=2}^\infty \frac{(\Delta_1+\Delta_2+\Delta_3) \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-1)!} + \sum_{u=2}^\infty \frac{\Delta_4 \Gamma(\varphi+\varepsilon)}{\Gamma(\varphi(u-1)+\varphi+\varepsilon)(u-2)!} \end{split}$$

Using hypothesis (13), with $\varepsilon = \varphi + \varepsilon$, we obtain

$$\begin{split} \hbar_5(\varepsilon,\varphi;\mu) &\leq \sum_{u=3}^\infty \frac{\Delta_1}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u-2}(u-3)!} + \sum_{u=2}^\infty \frac{3\Delta_1+\Delta_2}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u-2}(u-2)!} \\ &+ \sum_{u=2}^\infty \frac{\Delta_1+\Delta_2+\Delta_3}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u-2}(u-1)!} + \sum_{u=2}^\infty \frac{\Delta_4}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u-2}u!} \\ &= \sum_{u=0}^\infty \frac{\Delta_1}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u+1}u!} + \sum_{u=0}^\infty \frac{3\Delta_1+\Delta_2}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u}u!} \\ &+ \sum_{u=0}^\infty \frac{\Delta_1+\Delta_2+\Delta_3}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u}(u+1)!} + \sum_{u=0}^\infty \frac{\Delta_4}{(\varphi+\varepsilon)(\varphi+\varepsilon+1)^{u}(u+2)!} \\ &= \frac{\Delta_1(\varphi+\varepsilon+1)}{\varphi+\varepsilon} e^{\frac{1}{\varphi+\varepsilon+1}} + \frac{3\Delta_1+\Delta_2}{\varphi+\varepsilon} e^{\frac{1}{\varphi+\varepsilon+1}} + \frac{(\Delta_1+\Delta_2+\Delta_3)(\varphi+\varepsilon+1)}{\varphi+\varepsilon} \left(e^{\frac{1}{\varphi+\varepsilon+1}}-1\right) \\ &+ \frac{\Delta_4(\varphi+\varepsilon+1)^2}{\varphi+\varepsilon} \left(e^{\frac{1}{\varphi+\varepsilon+1}}-\frac{1}{\varphi+\varepsilon+1}-1\right) \leq \mu. \end{split}$$

3. Corollaries

By specializing the parameters Δ_1 , Δ_2 , Δ_3 , Δ_4 and μ in our Theorems, we obtain many results studied by many authors. The following is an illustration:

Let $\Delta_1=0$, $\Delta_2=\lambda$, $\Delta_3=1-\lambda$, $\Delta_4=-\gamma$ and $\mu=1-\gamma$ in Theorem 1. We conclude the subsequent corollary, which is due to [28] in Theorem 1.

Corollary 1. If $0 \le \gamma < 1$ and $0 \le \lambda < 1$, then the Rabotnov function $\mathcal{R}_{\ell,\beta}(\tau)$ belongs to the class $\mathcal{O}_{1-\gamma}(0,\lambda,1-\lambda,-\gamma)$ if

$$\left(\frac{\lambda\beta^2}{(\ell+1)^2} + \frac{(2\lambda+1)\beta}{\ell+1} + 1 - \gamma\right) e^{\frac{\beta}{\ell+1}} \leq 2(1-\gamma).$$

Let $\Delta_1=\lambda$, $\Delta_2=1-\lambda$, $\Delta_3=-\gamma$, $\Delta_4=0$ and $\mu=1-\gamma$ in Theorem 1. We conclude the subsequent corollary, which is due to [28] in Theorem 2.

Corollary 2. If $0 \le \gamma < 1$ and $0 \le \lambda < 1$, then the Rabotnov function $\mathcal{R}_{\ell,\beta}(\tau)$ belongs to the class $\mathcal{O}_{1-\gamma}(\lambda, 1-\lambda, -\gamma, 0)$ if

$$\left(\frac{\lambda\beta^3}{(\ell+1)^3} + \frac{(5\lambda+1)\beta^2}{(\ell+1)^2} + \frac{(4\lambda-\gamma+3)\beta}{\ell+1} + 1 - \gamma\right)e^{\frac{\beta}{\ell+1}} \le 2(1-\gamma).$$

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Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = \beta + 1$, $\Delta_4 = -\lambda(\gamma + \beta)$ and $\mu = 1 - \gamma$ in Theorem 2. We conclude the subsequent corollary, which is due to [29] in Theorem 1.

Corollary 3. If $0 \le \gamma < 1$, $0 \le \lambda \le 1$ and $\beta \ge 0$, then the Poisson distribution function $\mathcal{K}(\nu, \tau)$ belongs to the class $\mathcal{O}_{1-\gamma}(0,0,\beta+1,-\lambda(\gamma+\beta))$ if

$$(\beta + 1)\nu + (\beta + 1 - \lambda(\gamma + \beta))(1 - e^{-\nu}) \le 1 - \gamma.$$

Let $\Delta_1 = 0$, $\Delta_2 = \beta + 1$, $\Delta_3 = -\lambda(\gamma + \beta)$, $\Delta_4 = 0$ and $\mu = 1 - \gamma$ in Theorem 2. We conclude the subsequent corollary, which is due to [29] in Theorem 2.

Corollary 4. If $0 \le \gamma < 1$, $0 \le \lambda \le 1$ and $\beta \ge 0$, then the Poisson distribution function $\mathcal{K}(\nu, \tau)$ belongs to the class $\mathcal{O}_{1-\gamma}(0, \beta+1, -\lambda(\gamma+\beta), 0)$ if

$$(\beta + 1)\nu^2 + (3\beta + 3 - \lambda(\gamma + \beta))\nu + (\beta + 1 - \lambda(\gamma + \beta))(1 - e^{-\nu}) \le 1 - \gamma.$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 1 - \alpha \lambda$, $\Delta_4 = -\alpha(1 - \lambda)$ and $\mu = \beta |b|$ in Theorem 3. We conclude the subsequent corollary, which is due to [30] in Theorem 5.

Corollary 5. If $0 < \beta \le 1$, $0 \le \lambda \le 1$ and $b \in \mathbb{C} - \{0\}$, then the Bessel function $\mathcal{L}_{p,b}(u,c;\tau)$ belongs to the class $\mathcal{O}_{1-\alpha}(0,0,1-\alpha\lambda,-\alpha(1-\lambda))$ if

$$(1-\alpha\lambda)\mathcal{L}_p'(1)+(1-\alpha)\mathcal{L}_p(1)\leq 2(1-\alpha).$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 1 - \alpha \lambda$, $\Delta_4 = -\alpha(1 - \lambda)$ and $\mu = 1 - \alpha$ in Theorem 3. We conclude the subsequent corollary, which is due to [30] in Theorem 7.

Corollary 6. *If* $0 < \beta \le 1$, $0 \le \lambda \le 1$ *and* $b \in \mathbb{C} - \{0\}$, then the Bessel function $\mathcal{L}_{p,b}(u,c;\tau)$ belongs to the class $\mathcal{O}_{1-\alpha}(0,1-\alpha\lambda,-\alpha(1-\lambda),0)$ if

$$(1 - \alpha \lambda) \mathcal{L}''_n(1) + (3 - \alpha(1 + 2\lambda)) \mathcal{L}'_n(1) + (1 - \alpha) \mathcal{L}_p(1) \le 2(1 - \alpha).$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 2$, $\Delta_4 = -\cos \sigma - v$ and $\mu = \cos \sigma - v$ in Theorem 4. We conclude the subsequent corollary, which is due to [31] in Theorem 1.

Corollary 7. If $|\sigma| < \frac{\pi}{2}$ and $0 \le v < 1$, then $W^{(1)}(\varphi, \varepsilon; \tau)$ in $\mathcal{O}_{\cos \sigma - v}(0, 0, 2, -\cos \sigma - v)$ if

$$\varepsilon(\cos\sigma-v)+(\varepsilon+1)(\cos\sigma+v-2)\Big(e^{\frac{1}{\varepsilon+1}}-1\Big)-2e^{\frac{1}{\varepsilon+1}}\geq 0.$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 2$, $\Delta_4 = -\cos \sigma - v$ and $\mu = \cos \sigma - v$ in Theorem 4. We conclude the subsequent corollary, which is due to [31] in Theorem 2.

Corollary 8. If $|\sigma| < \frac{\pi}{2}$ and $0 \le v < 1$, then $W^{(1)}(\varphi, \varepsilon; \tau)$ in $\mathcal{O}_{\cos \sigma - v}(0, 2, -\cos \sigma - v, 0)$ if

$$\begin{split} & \varepsilon(\varepsilon+1)(\cos\sigma-v) - 2e^{\frac{1}{\varepsilon+1}} + (\varepsilon+1)(\cos\sigma+v-6)e^{\frac{1}{\varepsilon+1}} \\ & + (\varepsilon+1)^2(\cos\sigma+v-2)\Big(e^{\frac{1}{\varepsilon+1}}-1\Big) \geq 0. \end{split}$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 2$, $\Delta_4 = -\cos \sigma - v$ and $\mu = \cos \sigma - v$ in Theorem 5. We conclude the subsequent corollary, which is due to [31] in Theorem 3.

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Corollary 9. If $|\sigma| < \frac{\pi}{2}$ and $0 \le v < 1$, then $W^{(2)}(\varphi, \varepsilon; \tau)$ in $\mathcal{O}_{\cos \sigma - v}(0, 0, 2, -\cos \sigma - v)$ if

$$(\varphi + \varepsilon)(\cos \sigma - v) + (\cos \sigma + v)(\varphi + \varepsilon + 1)^{2} \left(e^{\frac{1}{\varphi + \varepsilon + 1}} - 1\right) + (\varphi + \varepsilon + 1)\left(2 - 2e^{\frac{1}{\varphi + \varepsilon + 1}} - \cos \sigma - v\right) \ge 0.$$

Let $\Delta_1 = \Delta_2 = 0$, $\Delta_3 = 2$, $\Delta_4 = -\cos \sigma - v$ and $\mu = \cos \sigma - v$ in Theorem 5. We conclude the subsequent corollary, which is due to [31] in Theorem 4.

Corollary 10. If $|\sigma| < \frac{\pi}{2}$ and $0 \le v < 1$, then $W^{(2)}(\varphi, \varepsilon; \tau)$ in $\mathcal{O}_{\cos \sigma - v}(0, 2, -\cos \sigma - v, 0)$ if

$$(\varphi+\varepsilon)(\cos\sigma-v)+(\cos\sigma+v-2)(\varphi+\varepsilon+1)\bigg(e^{\frac{1}{\varphi+\varepsilon+1}}-1\bigg)-2e^{\frac{1}{\varphi+\varepsilon+1}}\geq 0.$$

4. Conclusions

In this investigation, we determine sufficient conditions for the functions $\mathcal{R}_{\ell,\beta}(\tau)$, $\mathcal{K}(\nu,\tau)$, $\mathcal{L}_{p,b}(u,c;\tau)$, $\mathcal{W}^{(1)}(\varphi,\varepsilon;\tau)$ and $\mathcal{W}^{(2)}(\varphi,\varepsilon;\tau)$ through the general comprehensive subclass $\mathcal{O}_{\mu}(\Delta_1,\Delta_2,\Delta_3,\Delta_4)$ of univalent functions introduced in Definition 1. Furthermore, some corollaries of the results are also discussed. By suitably specializing the real constants $\Delta_1,\Delta_2,\Delta_3,\Delta_4$ and μ in our main results, we can determine new sufficient conditions for the functions $\mathcal{R}_{\ell,\beta}(\tau)$, $\mathcal{K}(\nu,\tau)$, $\mathcal{L}_{p,b}(u,c;\tau)$, $\mathcal{W}^{(1)}(\varphi,\varepsilon;\tau)$ and $\mathcal{W}^{(2)}(\varphi,\varepsilon;\tau)$ in other subclasses of analytic functions, which are new and have not been studied so far.

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